

超弦の場の理論における "多重Half-brane" Solution について

(On "Multiple Half-brane" Solutions
in Modified Cubic String Field Theory)

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0. Introduction

0. Introduction

String Field Theory

Non-perturbative string theory

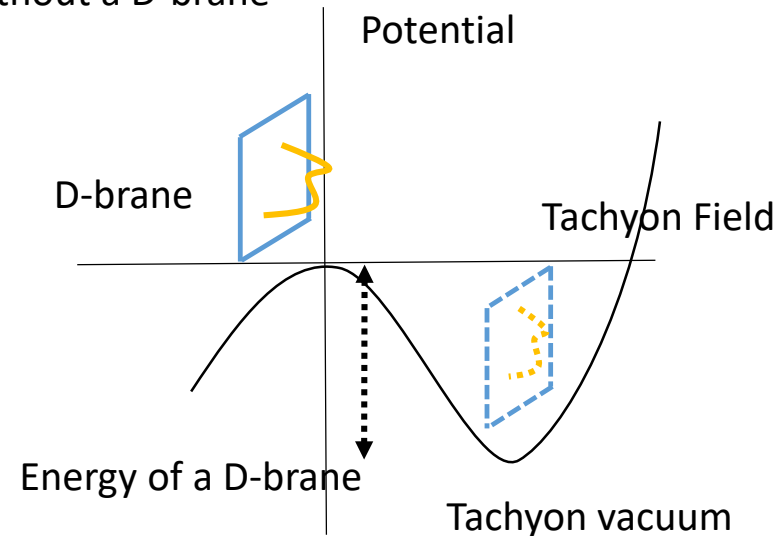
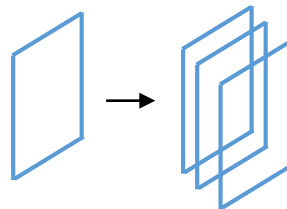
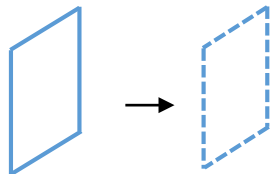
Sen's conjecture

Tachyon vacuum = the solution of string field theory without a D-brane

of D-brane $\rightarrow -1$



of D-brane $\rightarrow +n$??



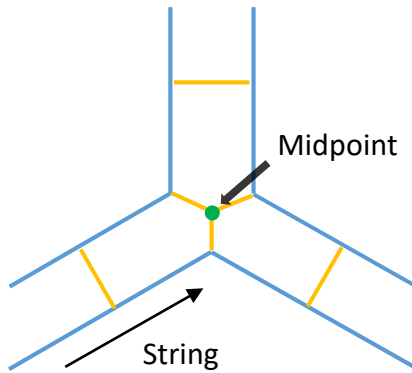
1. Bosonic Multiple Solution

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Action

$$- \int \left[\frac{1}{2} \Psi * Q\Psi + \frac{1}{3} \Psi * \Psi * \Psi \right]$$

[Witten '86]



Ψ	String Field : Open String, $gh\# = 1$, grassmann odd, defined by a superposition of all string states $\Psi = T(X)\hat{c}_1 0\rangle + A_\mu(X)\hat{\alpha}_{-1}^\mu\hat{c}_1 0\rangle + \frac{i}{\sqrt{2}}B(X)\hat{c}_0 0\rangle + \dots$
\int	BPZ inner product : String Field \rightarrow $c\#$
Q	BRS operator : String Field \rightarrow String Field
$*$	Star product (Midpoint int.): String Field $*$ String Field \rightarrow String Field

EOM

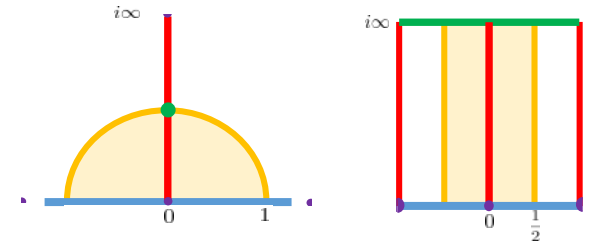
$$Q\Psi + \Psi^2 = 0$$

Gauge transformation

$$\delta_\Lambda \varphi = Q\Lambda + [\varphi, \Lambda]_\pm$$

(finite : $\varphi' = U^{-1}(Q + \varphi)U$)

1. Bosonic Multiple Solution



Sliver frame

Convenient coordinate of CFT to describe the $*$ product

[Rastelli-Zwiebach '01]

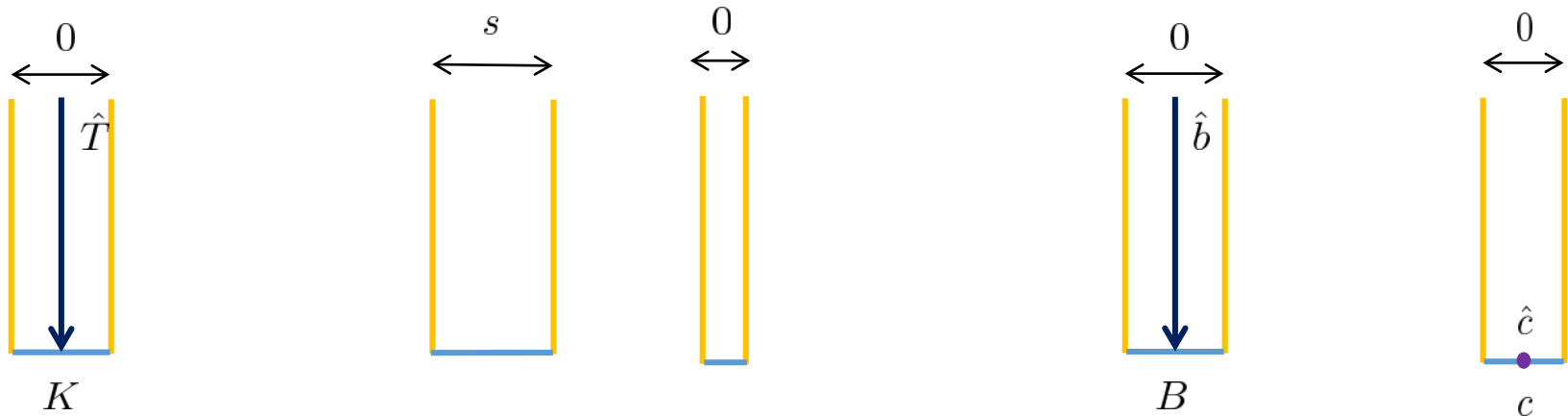
$$z = \frac{2}{\pi} \arctan w$$

w : UHP

String field K, B, c

K, B, c are defined by the sliver frame and used to construct solutions.

[Okawa '06]



$$K \rightarrow \int_{i\infty}^{-i\infty} \frac{dz}{2\pi i} \hat{T}(z)$$

$$e^{sK}$$

$$s \in \mathbb{R}_{\geq 0}$$

Wedge state

$$e^{0K} \rightarrow 1$$

Identity under the $*$ product

$$B \rightarrow \int_{i\infty}^{-i\infty} \frac{dz}{2\pi i} \hat{b}(z) |I\rangle$$

$$c \rightarrow \hat{c}(z)$$

\hat{b}, \hat{c} : ghost

\hat{T} : energy momentum tensor

1. Bosonic Multiple Solution

KBc alg.

closed under the following relations:

[Okawa '06]

[Nicholas-Schnabl '16]

$$\begin{aligned}
 [K, B] = 0 \quad \{B, c\} = 1 \quad B^2 = c^2 = 0 \\
 QB = K \quad QK = 0 \quad Qc = cKc
 \end{aligned}$$

BPZ inner product by using K, B, c in the sliver frame

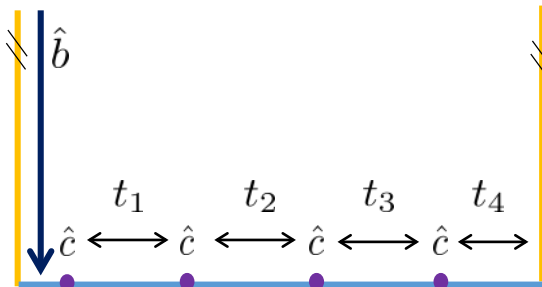
$$\int Bce^{t_1K} ce^{t_2K} ce^{t_3K} ce^{t_4K} = \int_{-i\infty}^{i\infty} \frac{dz}{2\pi i} \langle b(z)c(0)c(t_1)c(t_1+t_2)c(t_1+\dots) \rangle_{C_L} = -\frac{L^2}{4\pi^3} [t_3 \sin 2\theta_1 + \dots]$$

cf.) $\langle c(w_1)c(w_2)c(w_3) \rangle_{\text{UHP}} = w_{12}w_{13}w_{23}$

$(w_{ij} = w_i - w_j)$

$L = t_1 + \dots + t_4$

$\theta_i = \frac{\pi t_i}{L}$



1. Bosonic Multiple Solution

Pure-gauge form solution

[Okawa '06]

$$\Psi = U^{-1}QU \quad U = Bc + cBg(K)$$

* Formally, any solutions can be written in pure-gauge form by using the homotopy opr. .

$$U = 1 + A\Psi \quad QA = 1$$

[Ellwood '09]

Tachyon vacuum solution

NO D-brane = NO open string vacuum

[Schnabl '05
Erler-Schnabl '09]

$$\Psi_0 = U_1^{-1}QU_1 = -(Q(cB) + c)\frac{1}{1-K}$$

$$U_1 = Bc + cB\left(\frac{-K}{1-K}\right)$$

* we can treat $1/(1-K)$ as the superposition of wedge state by using the Schwinger parameter

$$\frac{1}{1-K} = \int_0^\infty dx e^{-x} e^{xK}$$

1. Bosonic Multiple Solution

Double-brane Solution

$$\Psi_2 = U_1 Q U_1^{-1} = -cB \frac{K^2}{1-K} c \frac{1}{K}$$

K_ϵ Regularization

$\frac{1}{K}$ is too singular to treat in the wedge state

$$\frac{1}{K} \rightarrow \frac{1}{K_\epsilon} = - \int_0^\infty dz e^{-\epsilon z} e^{zK}$$

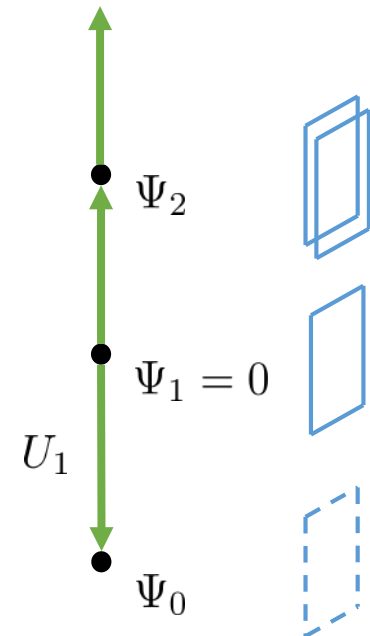
$$[\Psi_2]_\epsilon = -cB \frac{K_\epsilon^2}{1-K_\epsilon} c \frac{1}{K_\epsilon}$$

EOM in the strong sense

$$\int [[\Psi_2]_\epsilon (Q[\Psi_2]_\epsilon + [\Psi_2]_\epsilon^2)] = 0$$

$$\int [[\Psi_3]_\epsilon (Q[\Psi_3]_\epsilon + [\Psi_3]_\epsilon^2)] \neq 0$$

[Murata-Schnabl '11
Hata-Kojita '12]



$$\Psi_3 = U_1^2 Q U_1^{-2}$$

2. Half-brane Solution

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Modified cubic super SFT

[Preitschopf-Thorn-Yost
Arefeva-Medvedev-Zubarev '90]

$$S = - \int_{Y_{-2}} \left[\frac{1}{2} \Psi Q \Psi + \frac{1}{3} \Psi^3 \right]$$

Inverse PCO

$$\int_{Y_{-2}} \varphi = \langle \hat{Y}(i\infty) \hat{Y}(-i\infty) \hat{\varphi}_1(0) \rangle_{C_1}$$

GSO+ \rightarrow GSO+,- : “internal” CP factor

[Aref'eva-Belov-Giryavets '02]

$$\varphi \rightarrow \tilde{\varphi} \otimes \sigma_\mu$$

EOM

$$Q\Psi + \Psi^2 = 0$$

2. Half-brane Solution

$KBcG\gamma$ alg.

Extension of the KBc alg.

[Erler'11]

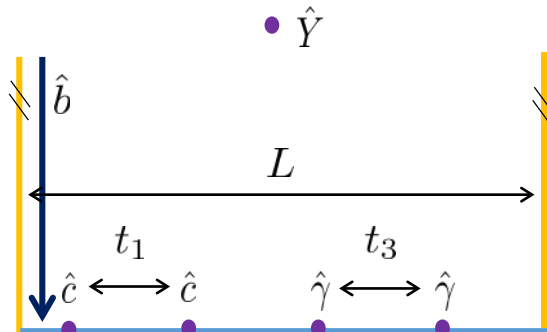
$$G \rightarrow \int_{-i\infty}^{i\infty} \frac{dz}{2\pi i} \hat{G}_{\text{supercurrent}}(z) \otimes \sigma_1 \quad \gamma \rightarrow \hat{\gamma}(z) \otimes \sigma_2$$

$$K \rightarrow \int_{-i\infty}^{i\infty} \frac{dz}{2\pi i} \hat{T}(z) \otimes \mathbb{I} \quad B \rightarrow \int_{-i\infty}^{i\infty} \frac{dz}{2\pi i} \hat{b}(z) \otimes \sigma_3 \quad c \rightarrow \hat{c}(z) \otimes \sigma_3$$

KBc alg +

$$\begin{aligned} \{G, G\} &= 2K, \quad [G, B] = 0, \quad [G, c] = 2\gamma, \quad \{G, \gamma\} = \frac{1}{2}\partial c \\ \{\gamma, B\} &= \{\gamma, c\} = 0 \\ QG &= 0, \quad Qc = cKc + \gamma^2, \quad Q\gamma = c\partial\gamma + \frac{1}{2}\gamma\partial c \end{aligned}$$

BPZ inner product by using K, B, c, γ in the sliver frame



$$\begin{aligned} \int_{Y_{-2}} B c e^{t_1} c e^{t_2} \gamma e^{t_3} \gamma e^{t_4} &= \left\langle \int \frac{dz}{2\pi i} \hat{Y}(i\infty) \hat{Y}(-i\infty) \hat{b}(z) \hat{c}(0) \hat{c}(t_1) \hat{\gamma}(t_1 + t_2) \hat{\gamma}(t_1 + \dots) \right\rangle_{C_L} \\ &= -\frac{t_1 L}{\pi^2} \cos\left[\frac{t_3 \pi}{L}\right] \end{aligned}$$

2. Half-brane Solution

Half-brane Solution

The solution constructed by using K, B, c, G, γ

[Erler '11]

$$\begin{aligned}\Psi_{\frac{1}{2}} &= U_{\frac{1}{2}}^{-1} Q U_{\frac{1}{2}} \\ &= -(Q(cB) + cBGc) \frac{1}{1-G}\end{aligned}$$

$$U_{\frac{1}{2}} = Bc + cB \frac{-G}{1-G}$$

$$\begin{aligned}\text{cf.) } U_1 &= Bc + cB \frac{-K}{1-K} \\ \Psi_0 &= -(Q(cB) + c) \frac{1}{1-K}\end{aligned}$$

The energy of this solution is half of the tachyon vacuum.

$$\frac{\pi^2}{3} \int_{Y_{-2}} [\Psi_{\frac{1}{2}} Q \Psi_{\frac{1}{2}}] = -\frac{1}{2}$$

* Since $1/K$ does not exist, we use the EOM.

2. Half-brane Solution

$$U_{\frac{1}{2}}^2 \text{ and } U_1$$

$$U_{\frac{1}{2}}^{-1} Q U_{\frac{1}{2}}^2 = (B\gamma^2 - cB(1-G)^2c) \frac{2G-1}{(1-G)^2} = \text{Tachyon Vacuum?}$$



$$\tilde{U} = U_{\frac{1}{2}}^{-2} U_1 = Bc + cB \frac{G-1}{1+G} = \text{regular gauge transformation?}$$



$$\int_{Y_{-2}} \tilde{\Psi} Q \tilde{\Psi} = 0 \dots ?? \rightarrow \text{Yes}$$

$$\tilde{\Psi} = \tilde{U}^{-1} Q \tilde{U} = -2(B\gamma^2 + cB \frac{1+G}{1-G} Kc) \frac{1}{1+G}$$

$$U_{\frac{1}{2}}^{-2} Q U_{\frac{1}{2}}^2 \text{ and } U_1^{-1} Q U_1 = \Psi_0 \text{ are gauge equivalent.}$$

2. Half-brane Solution

Classification by the sliver frame level expansion in [Erler '11]

$$U = Bc + cBg(K, G) \quad g(K, G) = g_+(K) + g_-(K)G$$

$$\text{Perturbative Vacuum} : \quad g_+(0) \neq 0$$

$$\text{Half Solution} : \quad g_+(0) = 0, \quad g_-(0) \neq 0$$

$$\text{Tachyon Vacuum} : \quad g_+(0) = 0, \quad g_-(0) = 0, \quad g'_+(0) \neq 0$$

$$U_{\text{half}} = Bc + cB[g_{h+}(K) + g_{h-}(K)G]$$

$$U_{\text{half}}^2 = Bc + cB[g_{h+}^2(K) + g_{h-}^2(K)K + 2g_{h+}(K)g_{h-}(K)G]$$

$$(g_{h+}^2(K) + g_{h-}^2(K)K)|_{K=0} = 0$$

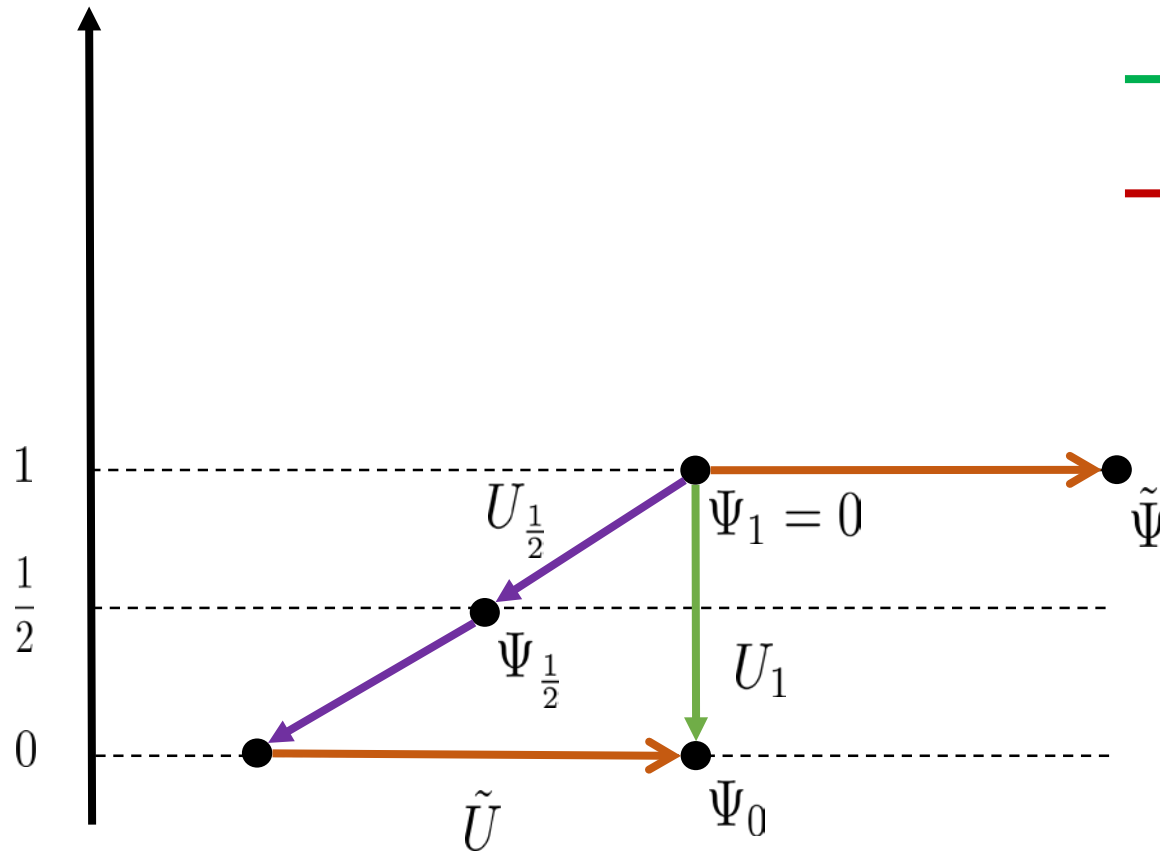
$$(2g_{h+}(K)g_{h-}(K))|_{K=0} = 0$$

$$(g_{h+}^2(K) + g_{h-}^2(K)K)'|_{K=0} = (g_{h-}^2(K))|_{K=0} \neq 0$$

$$\Rightarrow U_{\text{half}}^2 \sim U_{\text{tachyon vacuum}}$$

2. Half-brane Solution

$$1 + \frac{-\pi^2}{3} \int_{Y_{-2}} (U^{-1}QU)^3$$



● : Solution $U^{-1}QU$

→ $U_{\frac{1}{2}} = Bc + cB \frac{-G}{1-G}$

→ $U_1 = Bc + cB \frac{-K}{1-K}$

→ $\tilde{U} = Bc + cB \frac{G-1}{1+G}$



Perturbative vacuum



Tachyon vacuum

3. Double-brane Solution in Modified Cubic

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Double-brane Solution

$$\Psi_2 = U_1 Q U_1^{-1} = (B\gamma^2 + cB \frac{K^2}{1-K} c) \frac{1}{-K}$$

$$\text{cf.) } \Psi_0 = U_1^{-1} Q U_1$$

EOM in the strong sense

$$\int_{Y_{-2}} [\Psi_2]_\epsilon (Q[\Psi_2]_\epsilon + [\Psi_2]_\epsilon^2) = \frac{-2}{\pi^2 \epsilon} \neq 0$$

↔ In the **bosonic** case, the double-brane solution **satisfy** the EOM in the strong sense.

“winding number”

The double-brane solution done not satisfying the EOM in the strong sense, **but**

$$1 + \frac{-\pi^2}{3} \int_{Y_{-2}} [\Psi_2]_\epsilon^3 = 2$$

$$\text{Cf.) } 1 + \frac{-\pi^2}{3} \int_{Y_{-2}} \Psi_0^3 = 0$$

4. “Multiple Half-brane” Solution

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“Multiple Half-brane” Solution

$$\Psi? = U_{\frac{1}{2}} Q U_{\frac{1}{2}}^{-1} = -(-cBKc + B\gamma^2 + cBK \frac{1}{1-G} c) \frac{1}{G}$$

G_ϵ Regularization

$$(U_{\frac{1}{2}} = Bc + cB \frac{-G}{1-G})$$

$$G_\epsilon \equiv \tilde{G} \otimes \sigma_1 - \sqrt{-\epsilon} \otimes \sigma_3$$

$$\begin{aligned} \frac{1}{2} \{G_\epsilon, G_\epsilon\} &= \tilde{G}^2 \otimes \sigma_1^2 - \sqrt{-\epsilon} \tilde{G} \otimes \sigma_3 \sigma_1 - \sqrt{-\epsilon} \tilde{G} \otimes \sigma_1 \sigma_3 - \epsilon \otimes \sigma_3^2 \\ &= \tilde{G}^2 \otimes 1 - \epsilon \otimes 1 = \tilde{K} \otimes 1 - \epsilon \otimes 1 = K_\epsilon \end{aligned}$$

$$\Psi? \rightarrow [\Psi?]_\epsilon = -(-cBK_\epsilon c + B\gamma^2 + cBK_\epsilon \frac{1}{1-G_\epsilon} c) \frac{1}{G_\epsilon}$$

4. “Multiple Half-brane” Solution

EOM in the strong sense

$$\int_{Y_{-2}} [\Psi?]_{\epsilon} (Q[\Psi?]_{\epsilon} + [\Psi?]_{\epsilon}^2) = 0$$

“winding number”

$$1 + \frac{-\pi^2}{3} \int_{Y_{-2}} [\Psi?]_{\epsilon}^3 = \frac{3}{2}$$

$$\text{Cf.) } 1 + \frac{-\pi^2}{3} \int_{Y_{-2}} \Psi_{\frac{1}{2}}^3 = 1 + \frac{\pi^2}{3} \int_{Y_{-2}} \Psi_{\frac{1}{2}} Q \Psi_{\frac{1}{2}} = \frac{1}{2}$$

EOM in the strong sense + “winding number”

⇒ energy of this solution = (energy of a D-brane) × 3/2

4. "Multiple Half-brane" Solution

* Detail of the calculation

Ex.)
$$\int_{Y_{-2}} [\Psi?]_{\epsilon}^3 = -3 \int B\gamma^2 \frac{1}{G_{\epsilon}} cB \frac{K_{\epsilon}G_{\epsilon}}{1-G_{\epsilon}} c \frac{1}{G_{\epsilon}} cB \frac{K_{\epsilon}G_{\epsilon}}{1-G_{\epsilon}} c \frac{1}{G_{\epsilon}} + \dots$$

$$= 3 \int BcK_{\epsilon} \frac{1}{1-K_{\epsilon}} Kc\gamma K_{\epsilon} \frac{1}{1-K_{\epsilon}} \gamma \frac{1}{K_{\epsilon}} + \dots$$

$$= \dots$$

$$\int_{Y_{-2}} G_{\epsilon}\varphi = \frac{1}{2} \int_{Y_{-2}} \delta(\varphi)$$

(δ : super transformation)

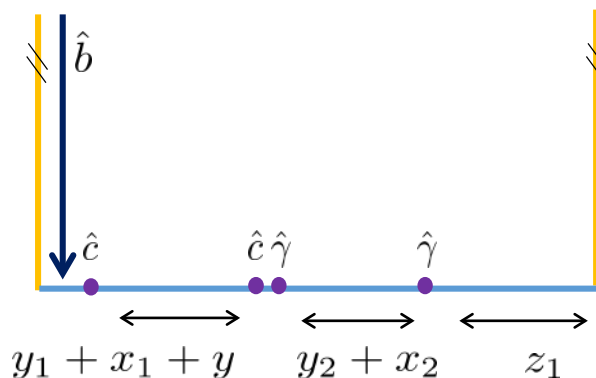
$$-\frac{\partial}{\partial y_1} \frac{\partial}{\partial y_2} \frac{\partial}{\partial y} \int_0^{\infty} dx_1 dx_2 dz e^{-x_1 - \epsilon - \epsilon(y_1 + y_2) - \epsilon z_1}$$

$$\left(\int_0^{\infty} da \int_0^1 db \int_0^1 dc a^2 \right.$$

$$a = x_1 + x_2 + z_1$$

$$b = \frac{z_1}{x_1 + x_2 + z_1}$$

$$c = \frac{z_1 + x_1}{x_1 + x_2 + z_1}$$



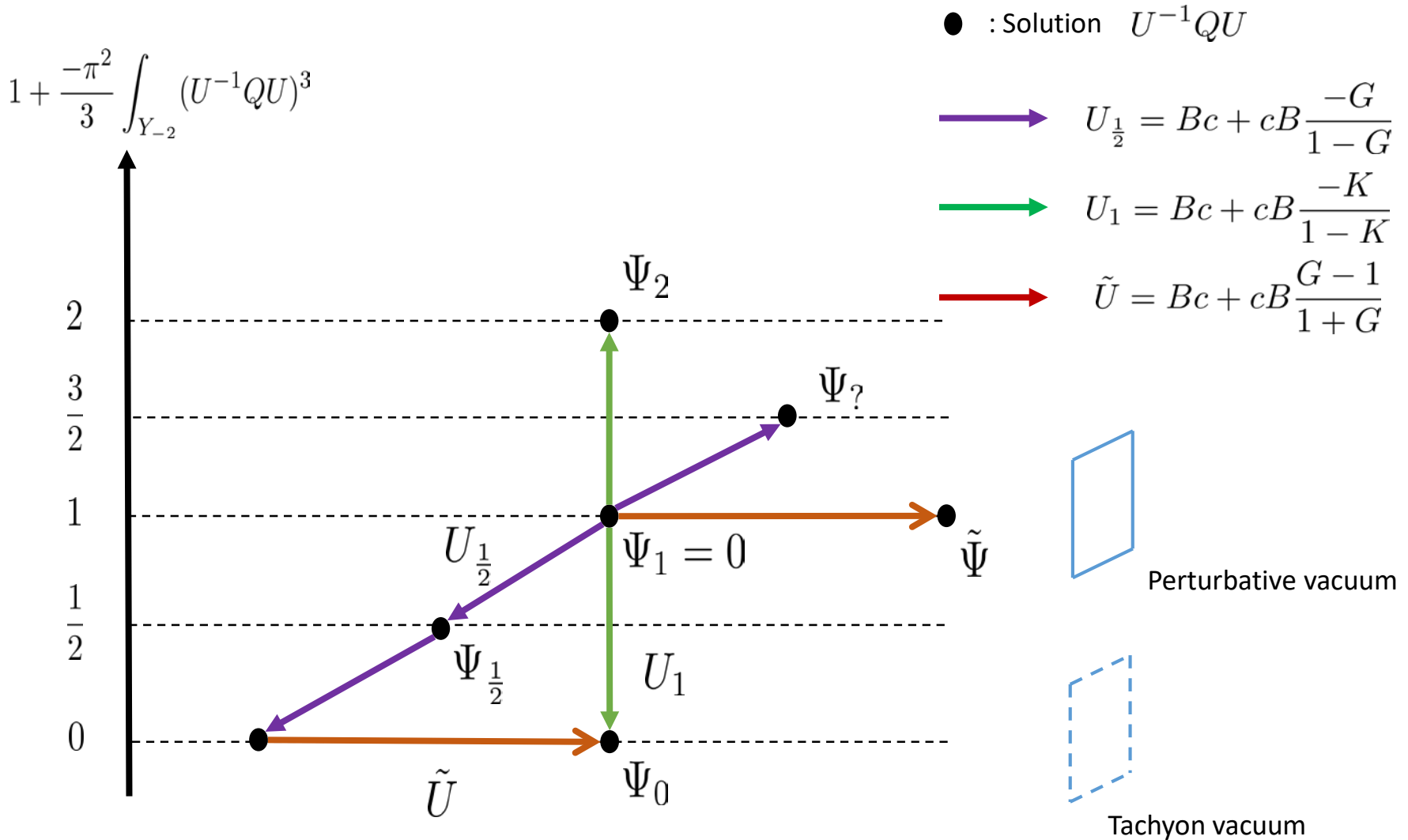
$$|_{y=0, y_1=0, y_2=0}$$

$$\parallel$$

$$-\frac{(y_1 + x_1 + y)L}{\pi^2} \cos\left[\frac{z_1\pi}{L}\right]$$

$$L = y_1 + x_1 + y + y_2 + x_2 + z_1$$

4. "Multiple Half-brane" Solution



5. Summary

5. Summary

Summary

As in bosonic case, we check **the EOM in the strong sense for the double-brane**, in modified cubic string field theory.

However, it does **not satisfy** the EOM and its energy is not twice of a D-brane. On the other hand, “winding number” is expected value.

We think the “**multiple half-brane**” solution constructed by the $KBcG\gamma$ alg. and we introduce the G_ϵ regularization, and we check that it satisfy the **EOM in the strong sense** and its energy is **(energy of a D-brane) $\times 3/2$** .

Future works

Multiple brane solution as the Erler-Maccaferri solution, which uses Boundary Condition Changing Operator in modified cubic string field theory

Solution of modified cubic string field theory
→ Solution of Berkovits' string field theory