

Flow equation, conformal symmetry and AdS geometries

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Ref.	S.Aoki-SY	ArXiv:1707.03982	PTEP (2018) 031B01
	S.Aoki-SY	ArXiv:1709.07281	NPB 933 (2018)
	S.Aoki-J.Balog-SY	ArXiv:1804.04636	

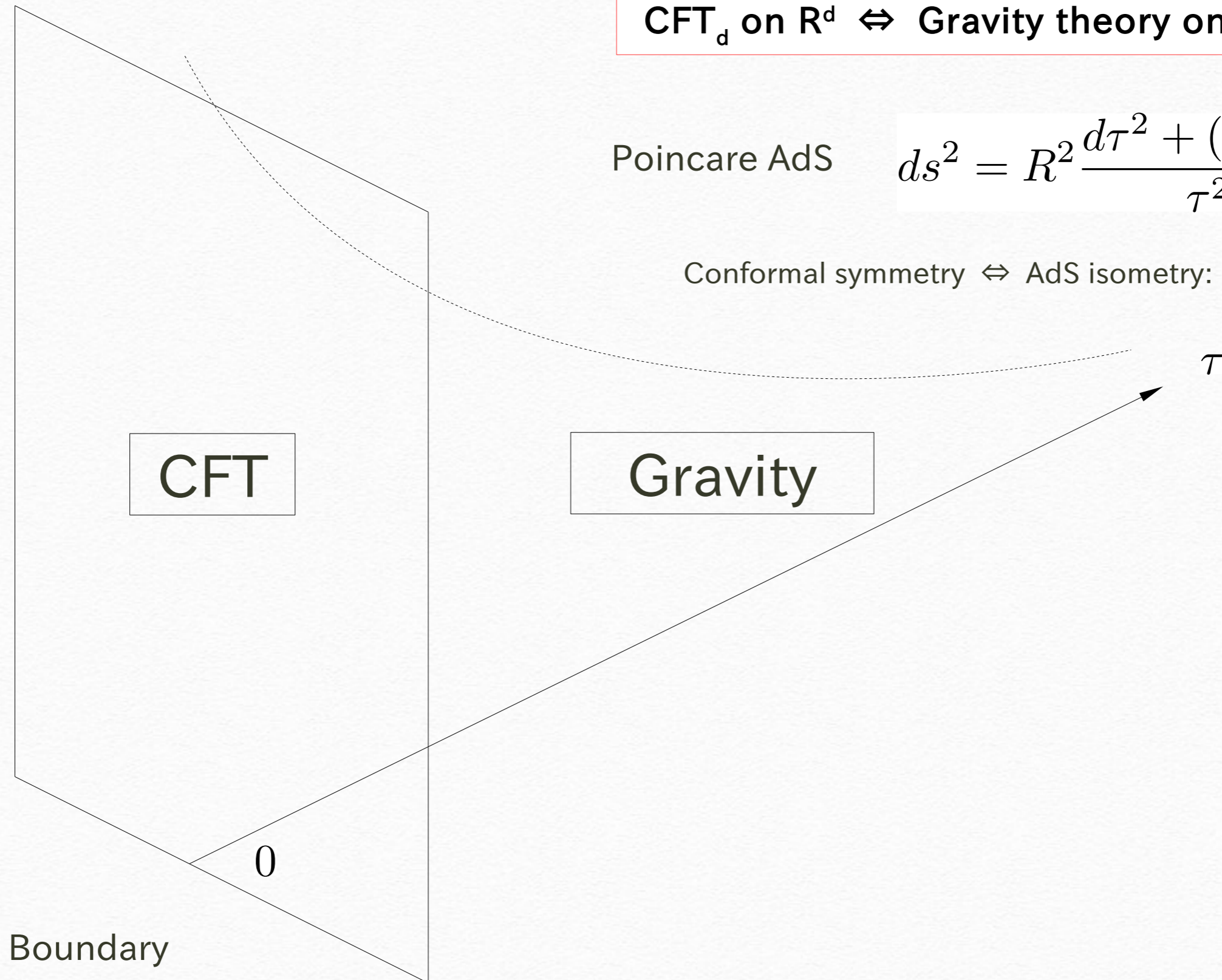
Holography and AdS/CFT

[Maldacena '97]

$$\text{CFT}_d \text{ on } \mathbb{R}^d \Leftrightarrow \text{Gravity theory on AdS}_{d+1}$$

Poincare AdS
$$ds^2 = R^2 \frac{d\tau^2 + (dx^\mu)^2}{\tau^2}$$

Conformal symmetry \Leftrightarrow AdS isometry: $SO(1, d+1)$



CFT

Gravity

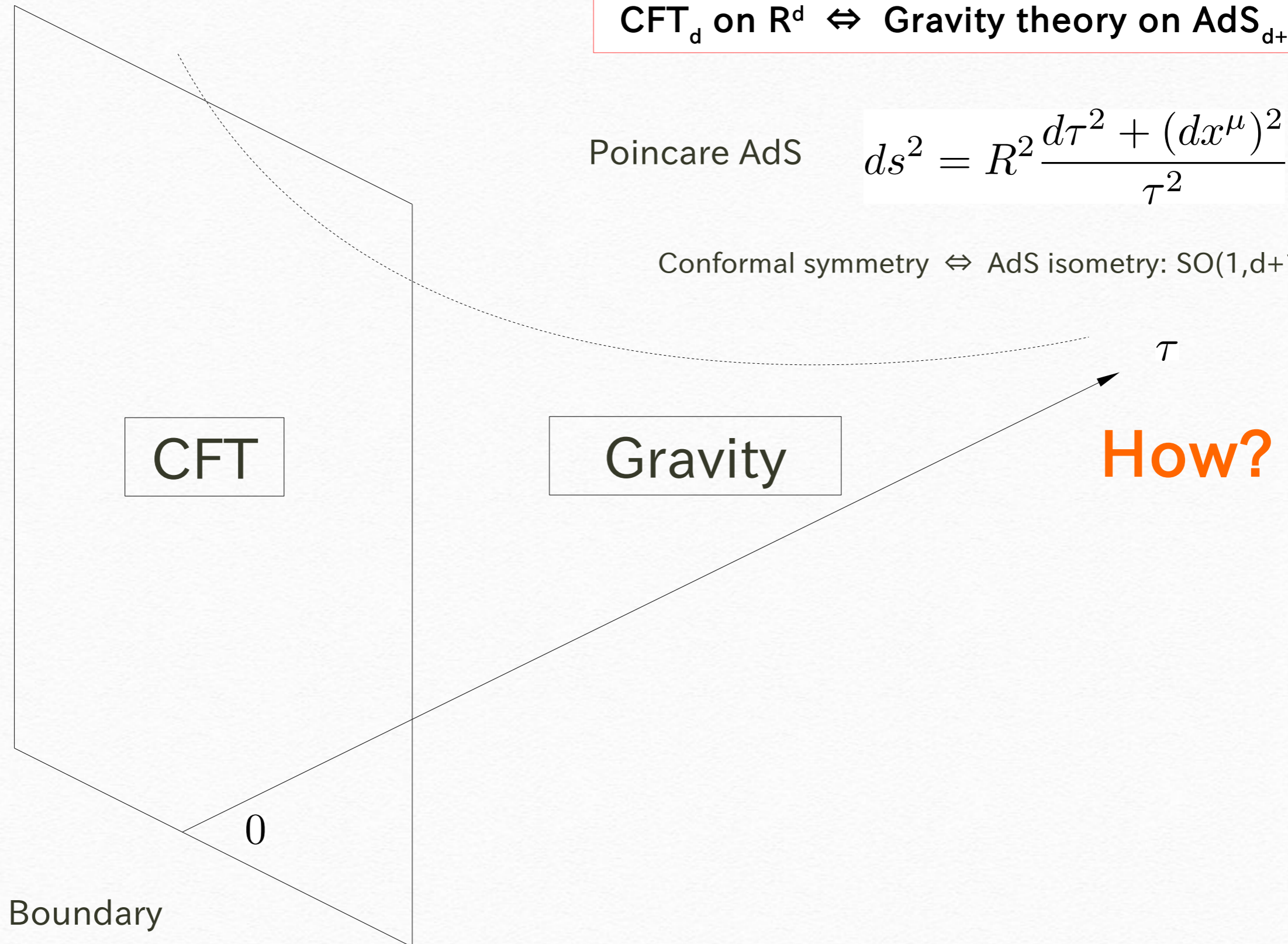
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AdS radial direction from CFT

1. Relevant RG flow

Construction of gravity solutions corresponding to UV and IR CFTs in the asymptotic regions.

[Girardello-Petrini-Porrati-Zaffaroni '98] [Distler-Zamora '98] [de-Boer-Verlinde-Verlinde '99] [Skenderis '00]

2. Wilsonian RG flow

The Wilsonian cut-off will correspond to sharp cut-off at the AdS radial direction

[Heemskerk-Polchinski '10]

3. Stochastic quantization

Euclidean path integral \equiv Equilibrium limit of statistical mechanical system coupled to a heat bath.

[Lifshytz-Periwal '00]

4. Entanglement entropy

continuous multi-scale entanglement renormalization ansatz (cMERA) [Swingle '09]...

5. bilocal field

Relative coordinate of bi-local field in vector models [Das-Jevicki '03]

6. Flow equation

Smearing operators so as to resolve a UV singularity in the coincidence limit

[Aoki-Kikuchi-Onogi '15] [Aoki-Balog-Onogi-Weisz '16,'17]

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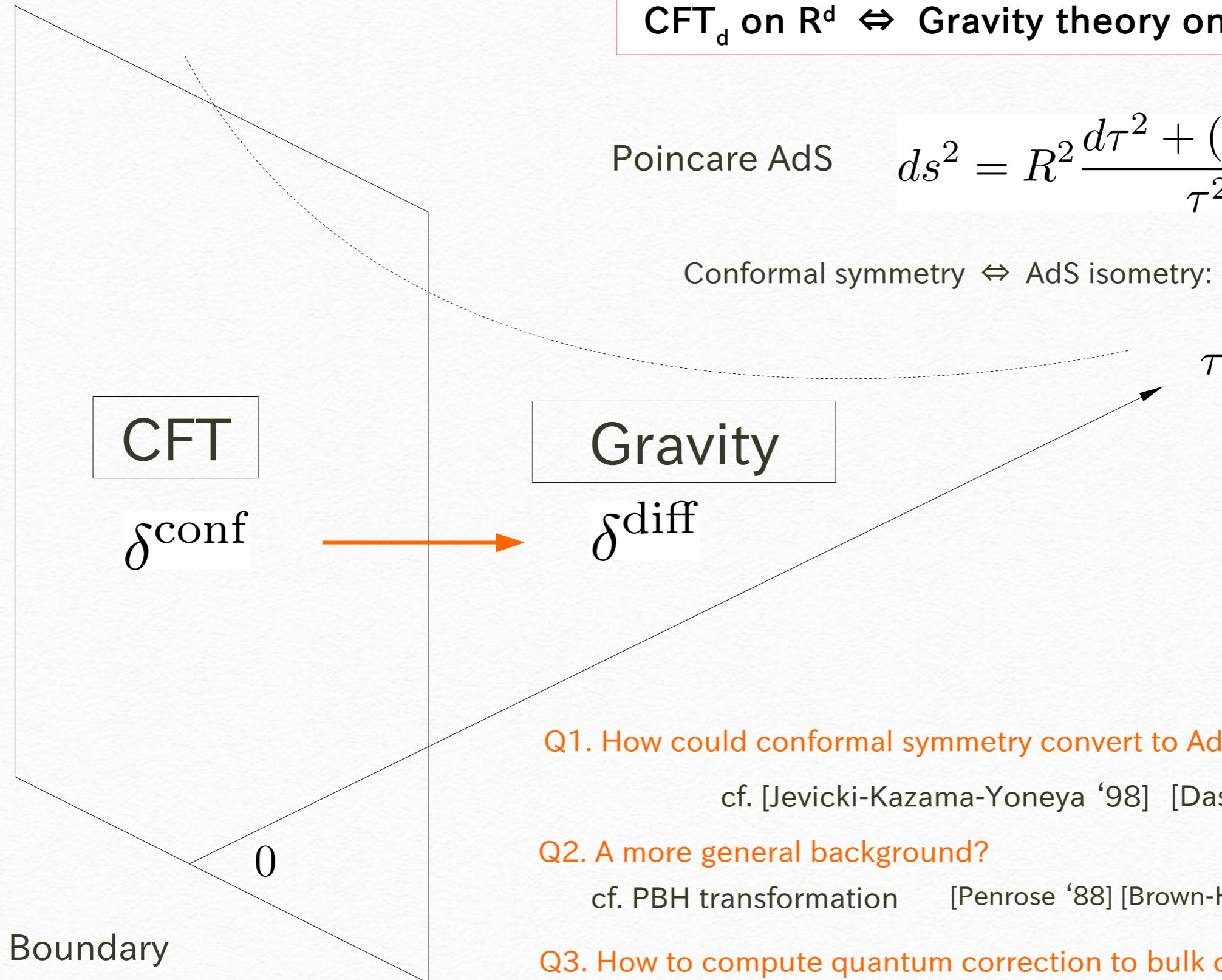
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Q1. How could conformal symmetry convert to AdS isometry?

cf. [Jevicki-Kazama-Yoneya '98] [Das-Jevicki '03]

Q2. A more general background?

cf. PBH transformation [Penrose '88] [Brown-Henneaux '86]

Q3. How to compute quantum correction to bulk observables?

Plan

✓ 1. Introduction

2. Flow equation & induced metric

Induced metric = information metric

3. Conformal symmetry \rightarrow AdS isometry

\leftarrow Answer for Q1

4. Generalization to conformally flat manifolds

4.1. Primary flow equation

\leftarrow Answer for Q2

4.2. AdS metric with conformally flat boundary

5. Summary

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Flow equation

(Gradient) flow equation

Equation to smear operators to resolve a UV singularity in the coincidence limit

[Albanese et al. (APE) '87] [Narayanan-Neuberger '06] [Luscher '10,'13]

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Consider QFT with n real component scalar field in d dimensions

Flow equation

$$\frac{\partial \phi^a(x; t)}{\partial t} = - \frac{\delta S_f(\varphi)}{\delta \varphi^a(x)} \Big|_{\varphi(x) \rightarrow \phi(x; t)} \quad \phi(x; 0) = \varphi(x)$$

t : flow time, $\phi(x; t)$: flowed field

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Comments

- (I) If S_f coincide with the action of the theory S
→ the flow equation is called the gradient one.
- (II) If S_f coincide with the free theory
→ the flow equation becomes the heat equation.
→ the operator ϕ is smeared similarly as in finite temperature t .
- (III) General solution → “flow kernel method”.

“Flow kernel method”

1. Rewrite the flow equation in the heat equation form.

$$\frac{\partial O(x; t)}{\partial t} = -\hat{H}O(x; t) \quad O(x; 0) = O(x)$$

2. Introduce the flow kernel (density).

$$\frac{\partial \rho(x, y; t)}{\partial t} = -\hat{H}\rho(x, y; t) \quad \rho(x, y; 0) = \delta^d(x - y)$$

3. A general solution is given by

$$O(x; t) = \int d^d y \rho(x, y; t) O(y)$$

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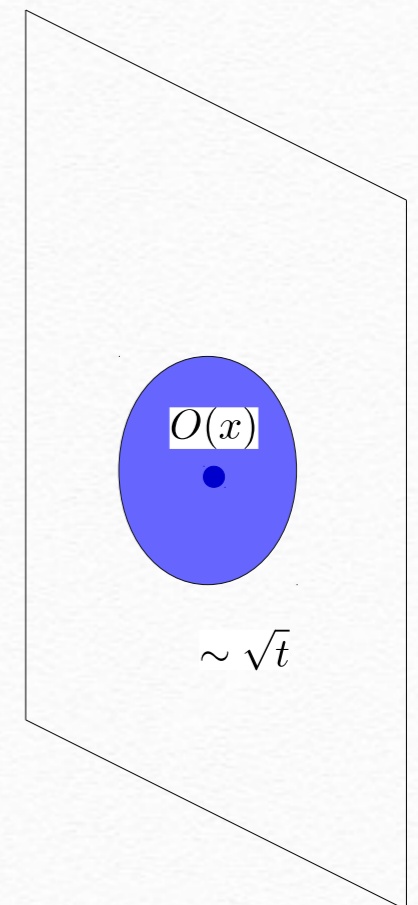
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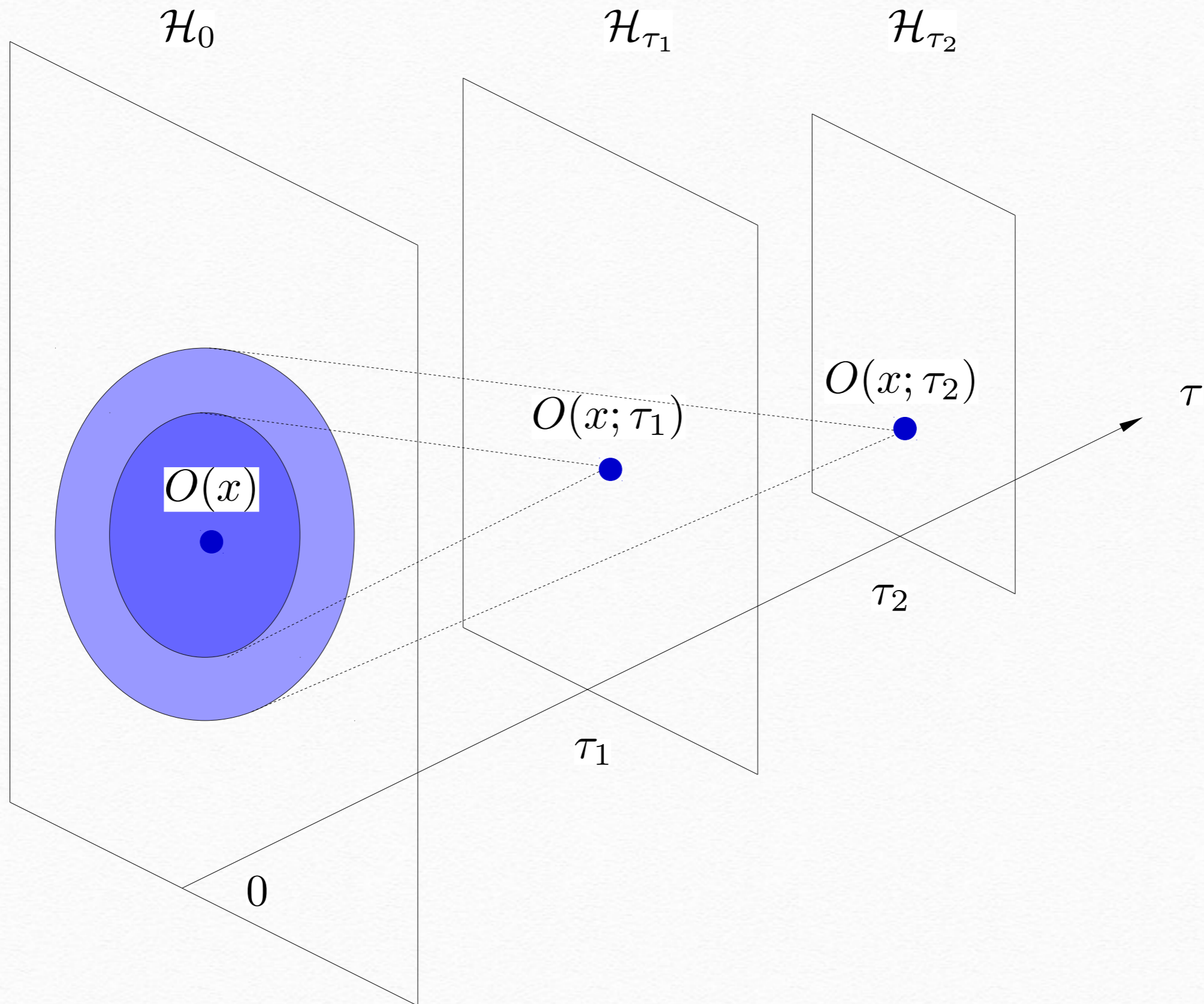
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Ex. Free flow

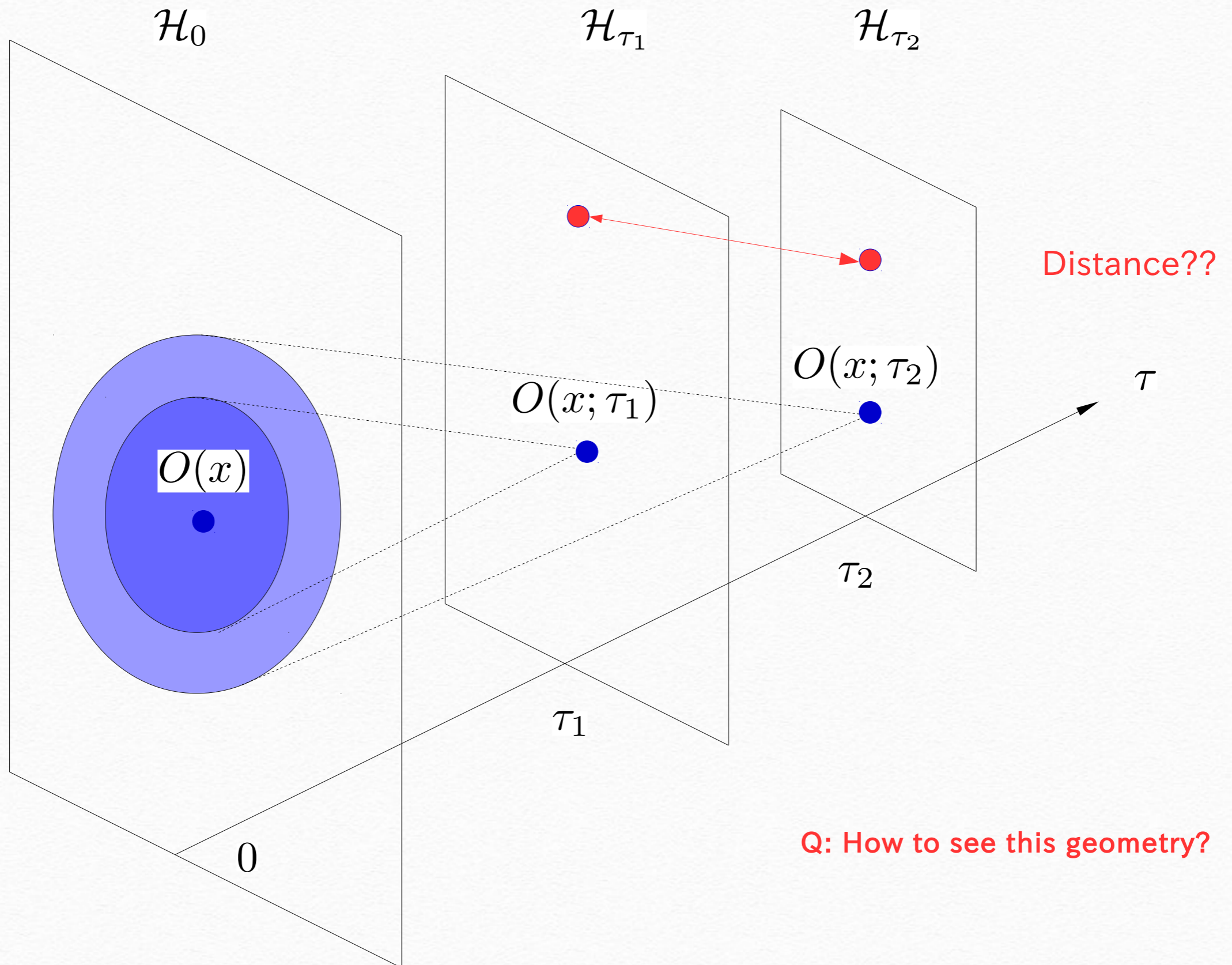
$$\rho(x, y; t) = \frac{1}{(4\pi t)^{d/2}} e^{-\frac{(x-y)^2}{4t}}$$



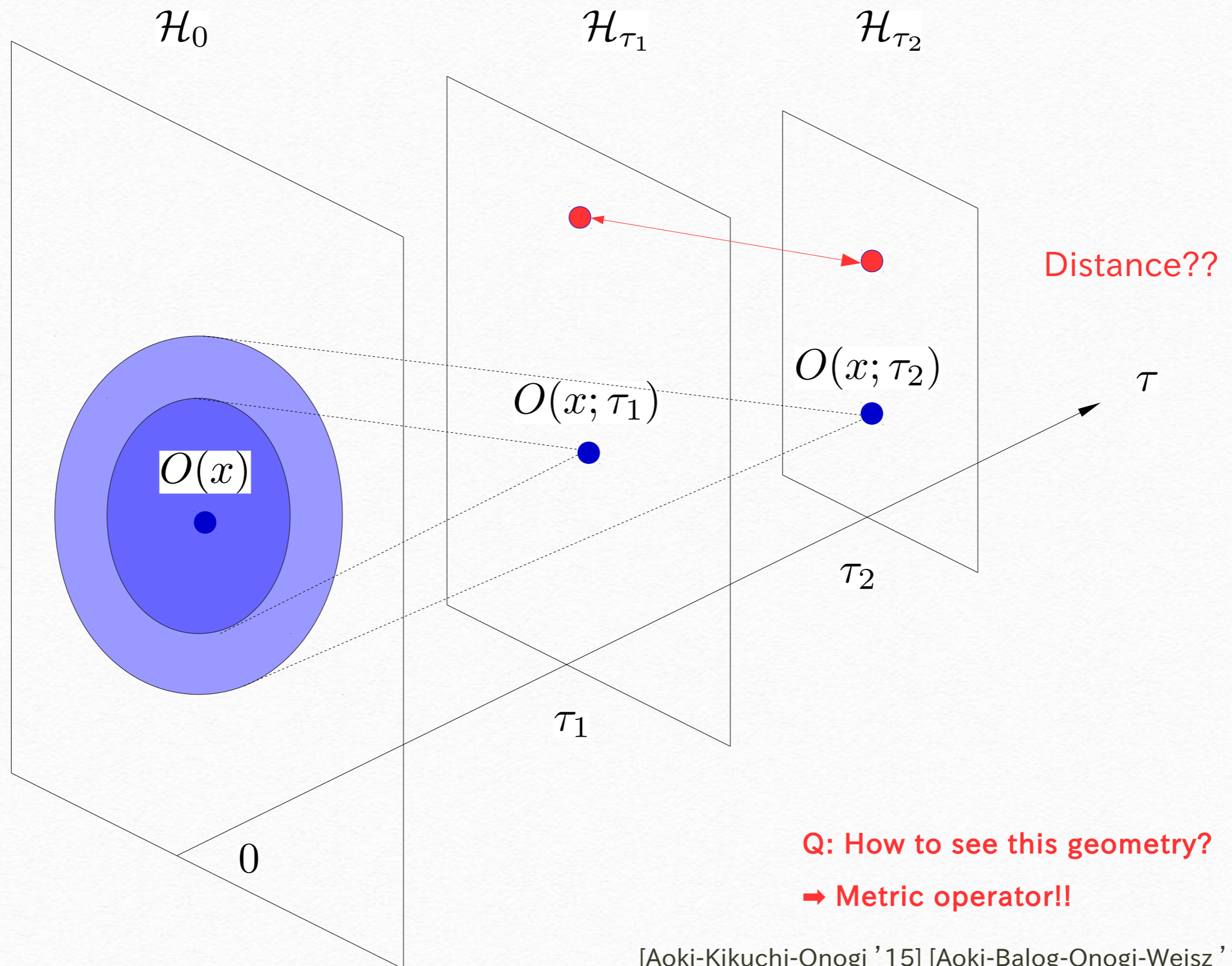
Sketch of smearing and extra direction



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Metric operator & induced metric

cf. [Aoki-Kikuchi-Onogi '15] [Aoki-Balog-Onogi-Weisz '16,'17]

Metric operator and induced metric

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Def. (Dimensionless normalized operator)

$$\sigma(x; t) := \frac{O(x; t)}{\sqrt{\langle O(x; t)^2 \rangle_{CFT}}}$$

“Operator renormalization”

NOTE: $\langle \sigma(x; t) \sigma(x; t) \rangle_{CFT} = 1$

Def. (Metric operator and induced metric)

$$\hat{g}_{MN}(x; t) := R^2 \frac{\partial \sigma(x; t)}{\partial z^M} \frac{\partial \sigma(x; t)}{\partial z^N} \quad g_{MN}(z) := \langle \hat{g}_{MN}(x; t) \rangle_{CFT}$$

R: constant of length dimension $z^M = (x^\mu, \tau)$ with $\tau = \sqrt{2dt}$

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Result 1 Induced metrics with the free flow equation become the AdS metric for a couple of vector models in the free or critical limit.

[Aoki-Kikuchi-Onogi '15] [Aoki-Balog-Onogi-Weisz '16,'17]

Result 2 An induced metric becomes quantum information metric for a general QFT.

[Aoki-SY '17]

Result 3 An induced metric defined in this way becomes AdS for a general CFT.

[Aoki-SY '17]

Induced metric = information metric

[S.Aoki-SY '17] PTEP (2018) 031B01

Induced metric = information metric

Def. (Bure metric for a density matrix)

$$D(\rho, \rho + d\rho)^2 = \frac{1}{2} \text{tr}(d\rho G)$$

ρ : density matrix G : hermitian 1 form operator satisfying $\rho G + G \rho = d\rho$

For a pure state, G is given by $G = d\rho$

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For $\forall z = (x, \tau)$, we can assign a density matrix such that

$$\rho_z = |\sigma(x; t)\rangle \langle \sigma(x; t)|, \quad |\sigma(x; t)\rangle := \sigma(x; t)|0\rangle, \quad \langle \sigma(x; t)| := \langle 0|\sigma(x; t)$$

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Def. (Inner product)

$$\langle \sigma(x; t) | \sigma(w; s) \rangle := \langle \sigma(x; t) \sigma(w; s) \rangle_{CFT}$$

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$$\langle\sigma(x; t)|\sigma(w; s)\rangle := \langle\sigma(x; t)\sigma(w; s)\rangle_{CFT}$$

We measure the infinitesimal distance between 2pt in the bulk manifold by the information metric for the density matrices of associated pure states.

$$ds_z^2 := D(\rho_z, (\rho + d\rho)_z)^2 = \frac{1}{2} \text{tr}(d\rho_z d\rho_z) = g_{MN}(z) dz^M dz^N,$$

cf. Fischer information metric $|\langle\Psi_\lambda|\Psi_{\lambda+\delta\lambda}\rangle|^2 = 1 - 2G_{\lambda\lambda}\delta\lambda^2 + \dots$

$|\Psi_\lambda\rangle$: Vacuum state for $H_0 + \lambda V$

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Correlator of free flowed field

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The 2pt function is $G_0(x; t|y; s) := \langle O(x; t)O(y; s) \rangle_{\text{CFT}} = e^{(t\partial_x^2 + s\partial_y^2)} \langle O(x)O(y) \rangle_{\text{CFT}}$

Poincare symmetry: $G_0(x; t|y; s) = \exists G_0((x - y)^2, t + s)$

Scaling property: $G_0(\lambda x; \lambda^2 t|\lambda y; \lambda^2 s) = \lambda^{-2\Delta} G_0(x; t|y; s)$

→ $G_0(x; t|y; s) = \frac{1}{(t + s)^\Delta} F_0\left(\frac{(x - y)^2}{t + s}\right)$ F_0 : a smooth function

Correlator of free flowed field

Consider a CFT with a scalar primary operator O with conformal dimension Δ

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The 2pt function of normalized field $G(x; t|y; s) := \langle \sigma(x; t)\sigma(y; s) \rangle_{\text{CFT}} = \left(\frac{2\sqrt{ts}}{t + s}\right)^\Delta F\left(\frac{(x - y)^2}{t + s}\right)$

$$F(x) \equiv F_0(x)/F_0(0), \quad F(0) = 1, \quad 2dF'(0) = -\Delta$$

Explicit computation: $F(x) = \frac{\Gamma(d/2)}{\Gamma(\Delta)\Gamma(d/2 - \Delta)} \int_0^1 dv v^{\Delta-1} (1 - v)^{d/2 - \Delta - 1} e^{-xv/4}$.

NOTE: UV singularity is resolved (regularized) $\Leftrightarrow F$ is a smooth function for $x > 0$.

Q1: Conformal symmetry \Rightarrow AdS isometry?

[S.Aoki-SY '17] PTEP (2018) 031B01

cf. [Jevicki-Kazama-Yoneya '98]

Conformal symmetry \rightarrow AdS isometry?

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Conformal transformation:

$$\delta x^\mu = a^\mu + \omega^\mu{}_\nu x^\nu + \lambda x^\mu + b^\mu x^2 - 2x^\mu (b_\nu x^\nu),$$
$$\delta^{\text{conf}} O(x) = -\delta x^\mu \partial_\mu O(x) - \frac{\Delta}{d} (\partial_\mu \delta x^\mu) O(x)$$

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\exists Higher derivative term \rightarrow Naively does not work for special conformal transformation...

cf. [Das-Jevicki '13] or Das's talk in 'Strings and Fields 2017'

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$$\delta^{\text{conf}} \hat{g}_{MN}(x; t) = \delta^{\text{diff}} \hat{g}_{MN}(x; t) + R^2 \lim_{(y; s) \rightarrow (x; t)} \frac{\partial}{\partial z^M} \frac{\partial}{\partial w^N} \{ \delta^{\text{extra}} \sigma(x; t) \sigma(y; s) + \sigma(x; t) \delta^{\text{extra}} \sigma(y; s) \}.$$

Conformal symmetry \rightarrow AdS isometry!!

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Vanish by taking VEV!!

$$\boxed{\langle \delta^{\text{conf}} \hat{g}_{MN}(x; t) \rangle = \langle \delta^{\text{diff}} \hat{g}_{MN}(x; t) \rangle .}$$

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Primary flow equation

How can we generalize the result of flat space boundary to a more general curved boundary?

A curved manifold need to admit CFT to live. → Restrict ourselves to a **conforamilly flat manifold**.

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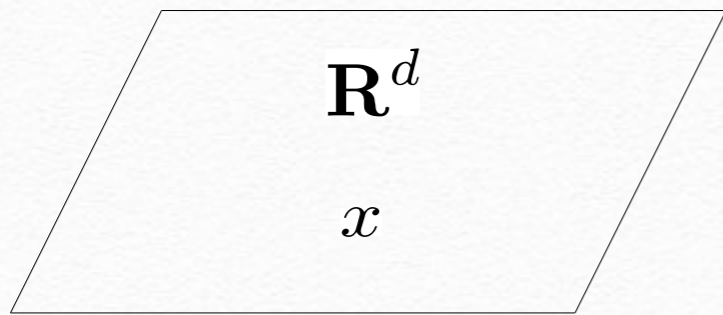
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We need to construct a flow equation associated with the curved manifold. But how?

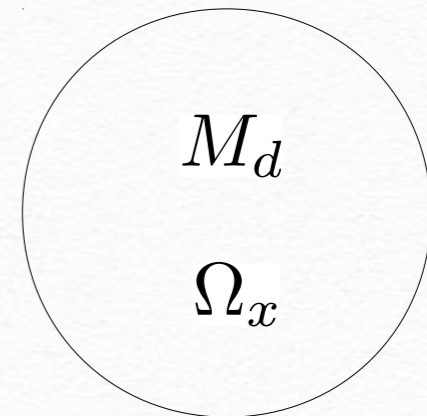
$$\frac{\partial O(x; t)}{\partial t} = \partial^2 O(x; t) \quad \rightarrow \quad ???$$

Setup



$$ds^2 = \delta_{\mu\nu} dx^\mu dx^\nu$$

Ω : conformal map



$$(ds^2)_{M_d} = g^{\frac{1}{d}}(x) \delta_{\mu\nu} dx^\mu dx^\nu$$

the conformal factor

$O(x)$

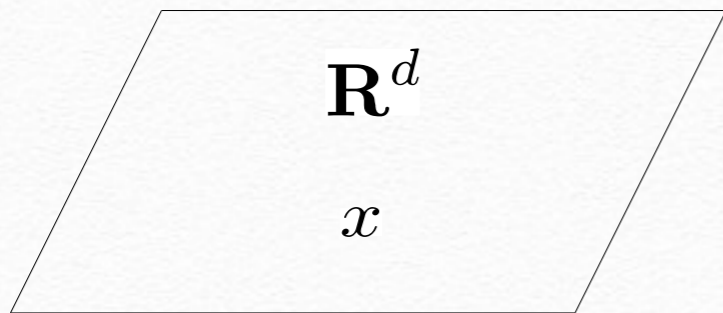
Ω^* : pullback



$O(\Omega_x)$

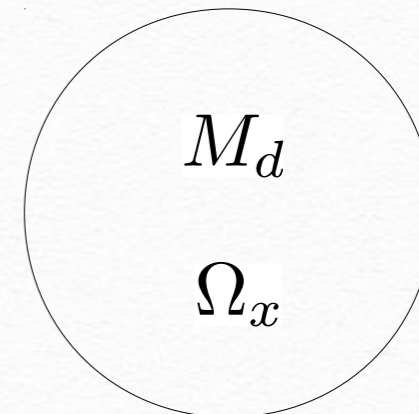
$$\Omega * O(\Omega_x) = U_\Omega^{-1} O(\Omega_x) U_\Omega = g^{-\frac{\Delta}{2d}}(x) O(x)$$

Primary flow



$$ds^2 = \delta_{\mu\nu} dx^\mu dx^\nu$$

Ω : conformal map
 \longrightarrow



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flow
 \downarrow

\downarrow flow

$O(x; t)$

???

$$\frac{\partial O(x; t)}{\partial t} = \partial^2 O(x; t)$$

Primary flow equation

How can we generalize the result of flat space boundary to a more general curved boundary?

A curved manifold need to admit CFT to live. → Restrict ourselves to a **conformally flat manifold**.

We need to construct a flow equation associated with the curved manifold. But how?

$$\frac{\partial O(x; t)}{\partial t} = \partial^2 O(x; t) \quad \rightarrow \quad ???$$

We request the following 2 conditions:

① There exists a flow time \tilde{t} associated with M_d corresponding to the flow time t on R^d such that the flowed operator inserted at Ω_x is related to the flowed one at x by the pullback of a conformal map Ω :

$$\Omega * O(\Omega_x; \tilde{t}) = U_\Omega^{-1} O(\Omega_x; \tilde{t}) U_\Omega = g^{-\frac{\Delta}{2d}}(x) O(x; t)$$

② The flow equation is invariant under the scale transformation.

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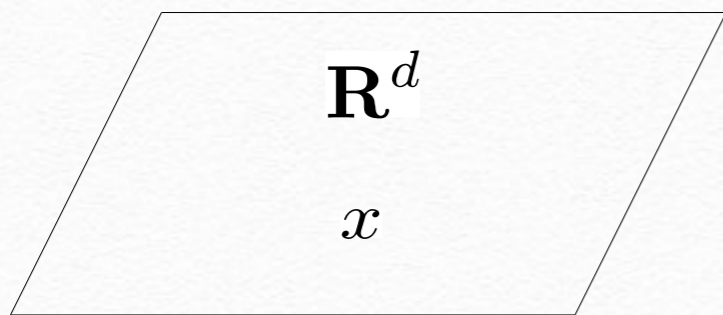
Primary flow equation



$$\tilde{t} = g^{\frac{1}{d}}(x)t$$

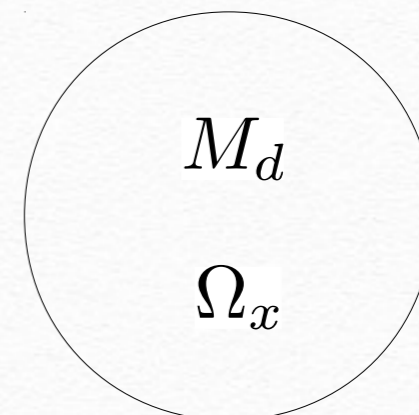
$$\frac{\partial}{\partial \tilde{t}} O(\Omega_x; \tilde{t}) = g^{-\frac{1}{d}}(x) g^{-\frac{\Delta}{2d}}(x) \partial^2 \cdot g^{\frac{\Delta}{2d}}(x) O(\Omega_x; \tilde{t}), \quad O(\Omega_x; 0) = O(\Omega_x)$$

Primary flow



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Ω : conformal map
 \longrightarrow



$$(ds^2)_{M_d} = g^{\frac{1}{d}}(x) \delta_{\mu\nu} dx^\mu dx^\nu$$

the conformal factor

$O(x)$

Ω^* : pullback
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flow
 \downarrow

$O(x; t)$

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Primary flow equation

Q2: Non-trivial curved boundary?

[S.Aoki-SY '17] NPB 933 (2018)

Generalization to curved background

A curved manifold needs to admit CFT to live. → Restrict ourselves to a **conformally flat manifold**.

We need to construct a flow equation associated with the curved manifold.

“Canonical” free flow equation (Primary flow equation)

$$\frac{\partial}{\partial \tilde{t}} O(\Omega_x; \tilde{t}) = g^{-\frac{1}{d}}(x) g^{-\frac{\Delta}{2d}}(x) \partial^2 \cdot g^{\frac{\Delta}{2d}}(x) O(\Omega_x; \tilde{t}), \quad O(\Omega_x; 0) = O(\Omega_x)$$

$g^{\frac{1}{d}}(x)$ the conformal factor

$$\tilde{t} = g^{\frac{1}{d}}(x) t$$

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$$g^{\frac{1}{d}}(x) \quad \text{the conformal factor} \quad \tilde{t} = g^{\frac{1}{d}}(x)t$$



The induced metric:

$$\tilde{g}_{\tilde{\tau}\tilde{\tau}}(z) = R^2 \frac{\Delta}{\tilde{\tau}^2},$$
$$\tilde{g}_{\tilde{\tau}\mu}(z) = g_{\mu\tilde{\tau}}(z) = -R^2 \frac{\Delta}{\tilde{\tau}} \frac{\partial}{\partial x^\mu} \log\{g^{\frac{1}{2d}}(x)\},$$
$$\tilde{g}_{\mu\nu}(z) = R^2 \Delta \left[\frac{\partial}{\partial x^\mu} \log\{g^{\frac{1}{2d}}(x)\} \frac{\partial}{\partial x^\nu} \log\{g^{\frac{1}{2d}}(x)\} + \frac{\delta_{\mu\nu} g^{\frac{1}{d}}(x)}{\tilde{\tau}^2} \right],$$

This is the (local) AdS metric whose radius is $R\sqrt{\Delta}$ with the boundary M_d !!

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This is the (local) AdS metric whose radius is $R\sqrt{\Delta}$ with the boundary M_d !!

This AdS metric can be obtained from the usual Poincare AdS by a bulk diffeomorphism

$$ds^2(\tilde{t}, x) = ds_{\text{PAdS}}^2(t, x)|_{t=g^{-1/d}(x)\tilde{t}}$$

Plan

- ✓ 1. Introduction
- ✓ 2. Flow equation & induced metric
- ✓ 3. Conformal symmetry \rightarrow AdS isometry
- ✓ 4. Generalization to conformally flat manifolds
 - 4.1. Primary flow equation
 - 4.2. AdS metric with conformally flat boundary
- 5. Quantum corrections
- 6. Summary

Q3: How to compute quantum correction?

[S.Aoki-J.Balog-SY '18]

Pre-geometric operators

Def. Pre-geometric operators are defined by replacing the metric which appears in the definition of the corresponding (classical) geometric object with the metric operator.

$$\mathcal{O}[g] \rightarrow \mathcal{O}[\hat{g}] =: \hat{\mathcal{O}}$$

Ex.

$$\hat{\Gamma}_{LN}^M(x; t) = \frac{1}{2} \hat{g}^{MP}(x; t) (\hat{g}_{P\{N,L\}}(x; t) - \hat{g}_{NL,P}(x; t))$$

$$\hat{R}_{LP}^M{}^N(x; t) = \partial_{[L} \hat{\Gamma}_{P]N}^M(x; t) + \hat{\Gamma}_{[LQ}^M(x; t) \hat{\Gamma}_{P]N}^Q(x; t)$$

$$\hat{R}_{PN}(x; t) = \hat{R}_{MP}^M{}^N(x; t),$$

$$\hat{R}(x; t) = \hat{g}^{PN}(x; t) \hat{R}_{PN}(x; t),$$

$$\hat{G}_{MN}(x; t) = \hat{R}_{MN}(x; t) - \frac{1}{2} \hat{g}_{MN}(x; t) \hat{R}(x; t).$$

Bulk interpretation

Taking the expectation value of these operators for a CFT state gives the observable in the bulk evaluated at its corresponding geometry, which obeys Einstein equation:

$$\langle \hat{G}_{AB} \rangle_\psi = T_{AB}^{\text{bulk}} |.$$

$$\langle \hat{G}_{AB} \rangle_\psi := \langle \psi | \hat{G}_{AB} | \psi \rangle$$

Ex. CFT vacuum \rightarrow Only cosmological constant.

$$T_{AB}^{\text{bulk}} | = -\Lambda g_{AB}$$

CLAIM: This induced Einstein tensor is expected to describe **that of dual quantum gravity**.

In particular, let us compute LHS in the **1/n expansion**:

$$\langle G_{AB}[\hat{g}] \rangle = G_{AB}[\langle \hat{g} \rangle] + \langle G_{AB}[\hat{g}] \rangle_c$$

Leading Order (LO)

Next to Leading Order (NLO)

Classical geometry

Quantum correction

Quantum correction to bulk CC

Ex. CFT vacuum, $O(n)$ free vector model \rightarrow Quantum correction to cosmological constant.

① Expand the operator around the vacuum:

$$\hat{G}_{AB} = G_{AB} + \dot{G}_{AB} + \ddot{G}_{AB} + \dots$$

$$\mathcal{O}(h^0) \quad \mathcal{O}(h^1) \quad \mathcal{O}(h^2) \quad h_{AB} = \hat{g}_{AB} - g_{AB}$$

$$G_{AB} = G_{AB}[\langle \hat{g} \rangle] = \frac{d(d-1)}{2L^2\Delta} g_{AB}$$

$$\Lambda = -\frac{d(d-1)}{2L^2\Delta} + \mathcal{O}\left(\frac{1}{n}\right),$$

$$\ddot{G}_{AB} = R_{AB} - \frac{1}{2}h_{AB}\dot{R} - \frac{1}{2}g_{AB}\ddot{R}$$

② Taking the VEV:

$$\langle \hat{G}_{AB} \rangle = G_{AB} + \langle \ddot{G}_{AB} \rangle + \dots$$

$$\langle \ddot{G}_{AB} \rangle = \frac{(d-1)d(d+4)}{2nL^2\Delta} g_{AB}.$$

$$\Lambda = -\frac{d(d-1)}{2L^2\Delta} \left(1 + \frac{d+4}{2n}\right) + \mathcal{O}\left(\frac{1}{n^2}\right), \quad L_{\text{AdS}}^2 = L^2\Delta \left(1 - \frac{d+4}{2n}\right) + \mathcal{O}\left(\frac{1}{n^2}\right).$$

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Summary

- An induced metric defined from the normalized flowed field generally corresponds to the **quantum information metric**.
- For a general CFT, we explicitly showed that the induced metric with the free flow equation always becomes the **AdS** metric.
- Conformal symmetry converts to AdS isometry **after quantum averaging**.
- We have constructed a canonical flow equation for a primary scalar operator on a conformally flat manifold, called the **primary flow equation**.
- The induced metric associated with the primary flow equation becomes AdS with **the conformally flat boundary**.
- AdS with the conformally flat boundary is obtained from the usual Poincare AdS **by a simple bulk diffeomorphism transformation**.
- Demonstrate how to compute **quantum corrections to bulk observables** and compute it explicitly for the **cosmological constant**.

Future works

- How to encode dynamics beyond geometry? For excited states?

working in progress [Aoki-Balog-SY]

- How to reconstruct bulk operator?

- Canonical choice of a normalized field?

- Finite temperature? BH?

- Spin 1,2 field? Fermion?

- de Sitter construction? Application to the real world?

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Thank you!